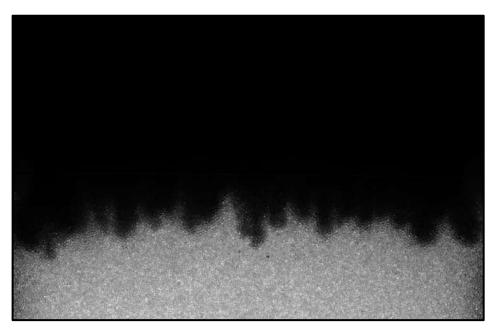
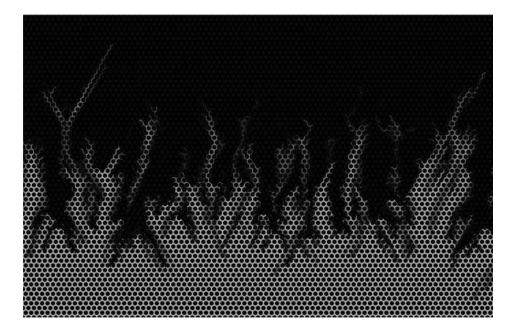
Pore-scale simulation of convective mixing in confined media





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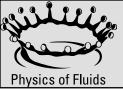


⁵Max Plank Institute for Dynamics and Self-Organization, Göttingen (Germany)

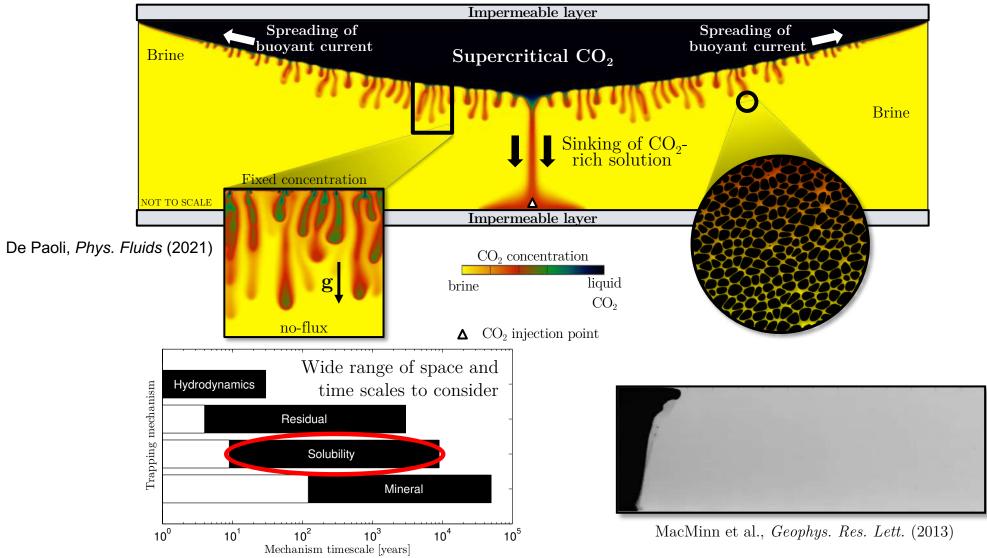
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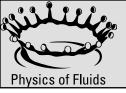
September 10-13, 2023 Vienna (Austria)



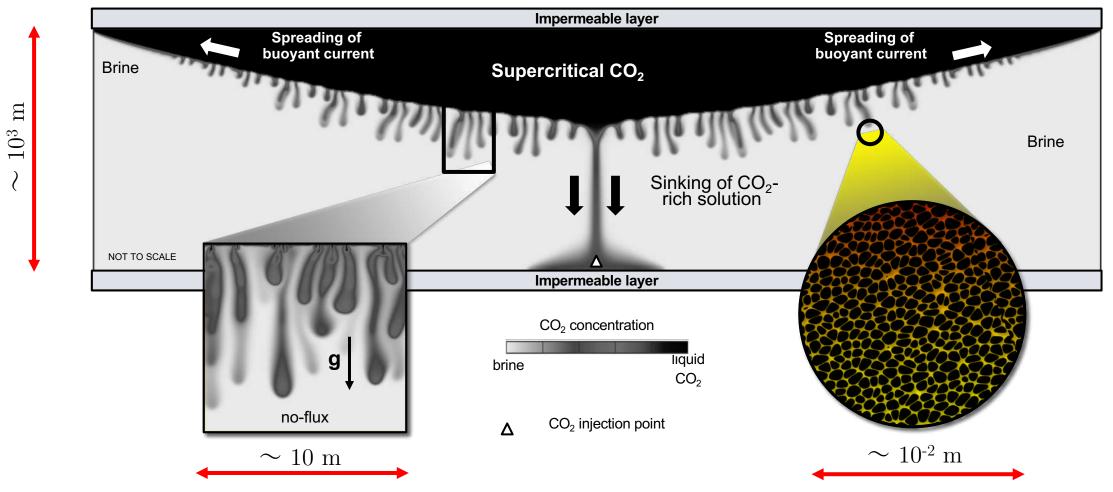
Carbon Capture and Storage





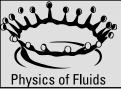


Convection in complex multiphase and multiscale systems

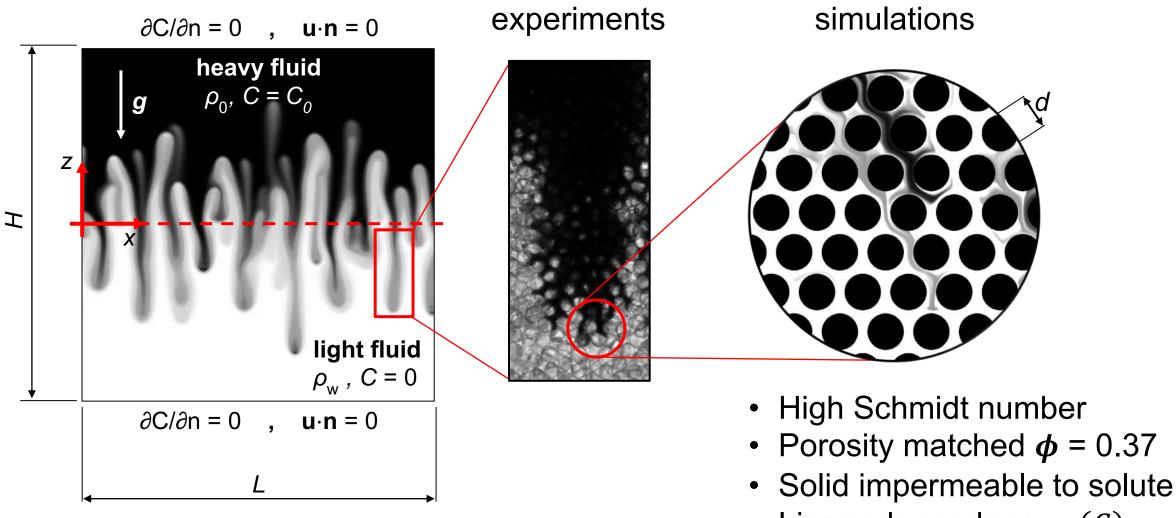


De Paoli, Phys. Fluids (2021)





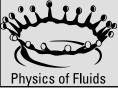
Flow configuration



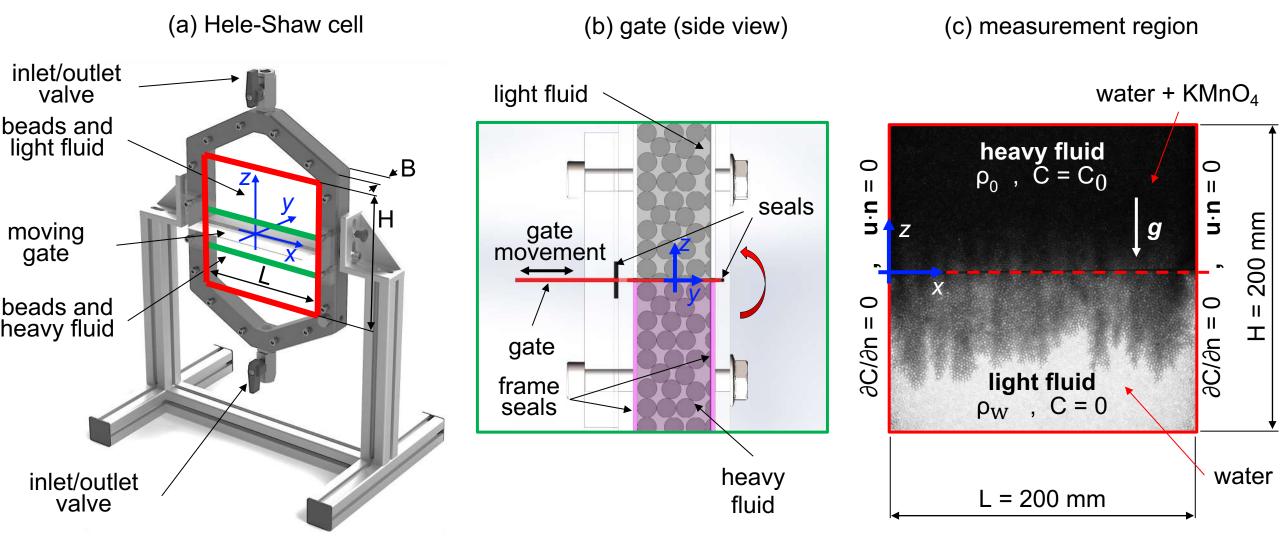


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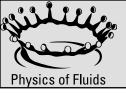
Computational Mechanics



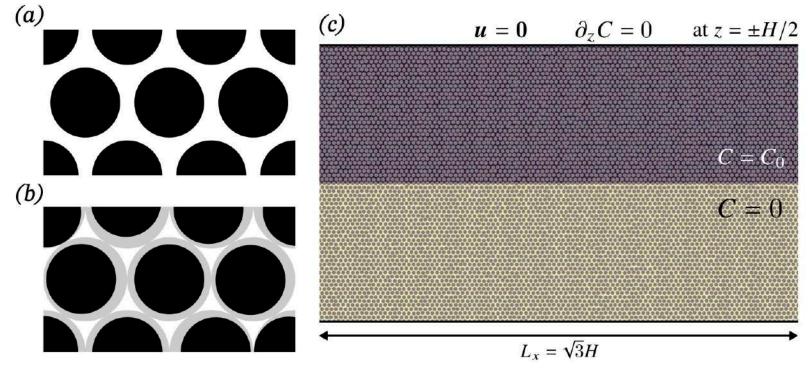
Experimental setup







Numerical method



$$\partial_t \boldsymbol{u} + (\boldsymbol{u} \cdot \boldsymbol{\nabla})\boldsymbol{u} = -\rho_0^{-1}\boldsymbol{\nabla}p + \boldsymbol{v}\nabla^2\boldsymbol{u} - g\beta C\hat{\boldsymbol{z}},$$

$$\partial_t C + (\boldsymbol{u} \cdot \boldsymbol{\nabla})C = D\nabla^2 C,$$

$$\rho = \rho_0 \left[1 + \frac{\Delta\rho}{\rho_0 C_0} (C - C_0)\right]$$

Pore-scale simulation of convective mixing in confined media *Marco De Paoli,* Physics of Fluids Group, University of Twente Advanced finite difference (AFiD, open source) + Immersed Boundaries Method

Resolution:

- velocity: ≥ 32 points per diameter
- concen.: ≥ 128
 points per diameter

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Computational Mechanics



withing.
Physics of Fluids

experiments											
Name	H/d	ϕ	Sc	Ra	<i>Ra</i> _d	Ra^*	Pe	Re			
				10	2	2					
E1	200	0.37	558	4.535×10^{10}	5.669×10^3	2.173×10^{3}	0.289	0.0005			
E2	200	0.37	558	9.099×10^{10}	1.137×10^{4}	4.359×10^{3}	0.580	0.0010			
E3	200	0.37	558	1.824×10^{11}	2.280×10^4	8.737×10^{3}	1.163	0.0021			
E4	200	0.37	558	3.637×10^{11}	4.546×10^4	1.742×10^{4}	2.320	0.0042			
E5	114	0.37	558	4.667×10^{10}	3.126×10 ⁴	6.846×10^3	1.595	0.0029			
E6	114	0.37	558	9.099×10^{10}	6.096×10^4	1.335×10^{4}	3.110	0.0056			
E7	114	0.37	558	1.820×10^{11}	1.219×10^{5}	2.671×10^4	6.222	0.0112			
E8	114	0.37	558	3.626×10^{11}	2.429×10^{5}	5.320×10^4	12.395	0.0222			
E9	67	0.35	558	4.490×10 ¹⁰	1.515×10^{5}	1.627×10^4	5.795	0.0104			
E10	67	0.35	558	9.495×10 ¹⁰	3.204×10^{5}	3.441×10^4	12.256	0.0220			
E11	67	0.35	558	1.834×10^{11}	6.189×10^5	6.646×10^4	23.672	0.0425			
E12	67	0.35	558	3.670×10^{11}	1.239×10^{6}	1.330×10^{5}	47.370	0.0850			
E13	50	0.37	558	4.506×10 ¹⁰	3.605×10 ⁵	3.454×10^4	18.393	0.0330			
E14	50	0.37	558	9.101×10^{10}	7.281×10^{5}	6.976×10^4	37.150	0.0666			
E15	50	0.37	558	1.824×10^{11}	1.460×10^{6}	1.398×10^{5}	74.474	0.1336			
E16	50	0.37	558	3.622×10 ¹¹	2.898×10^{6}	2.777×10^5	147.861	0.2652			

flow scales and parameters

 $k = \frac{d^2}{36k_C} \frac{\phi^3}{(1-\phi)^2}$

$$U = \frac{g \Delta \rho k}{\mu} \qquad \qquad \ell = \frac{\phi D}{U}$$

$$Sc = \frac{\mu}{\rho_0 D}$$

 $Da = k/H^2$

 $Ra^* = \frac{Ra\,Da}{\phi}$

simulations											
Name	H/d	ϕ	Sc	Ra	Ra _d	Ra*	Pe	Re			
S1 S2	17 17	0.37 0.37	100 100	5.268×10^{8} 1.666×10^{9} 5.268×10^{9}	1.000×10^{5} 3.162×10^{5} 1.000×10^{6}	3.334×10^{3} 1.054×10^{4} 3.334×10^{4}	5.102 16.135	0.0510 0.1614			
S3 S4 S5	17 35 35	0.37 0.37 0.37	100 100 100	3.208×10^{-10} 4.214×10^{9} 1.333×10^{10}	1.000×10^{5} 3.162×10^{5}	6.669×10^{3} 2.109×10^{4}	51.024 5.102 16.135	0.5102 0.0510 0.1614			
<u>S6</u>	35	0.37	100	4.214×10 ¹⁰	1.000×10^{6}	6.669×10 ⁴	51.024	0.5102			
S7 S8 S9	52 52 52	0.37 0.37 0.37	100 100 100	$\begin{array}{c} 1.422 \times 10^{10} \\ 4.498 \times 10^{10} \\ 1.422 \times 10^{11} \end{array}$	1.000×10^{5} 3.162×10^{5} 1.000×10^{6}	1.000×10^{4} 3.163×10^{4} 1.000×10^{5}	5.102 16.135 51.024	0.0510 0.1614 0.5102			
\$10 \$11 \$12	70 70 70	0.37 0.37 0.37	100 100 100	$\begin{array}{c} 3.372 \times 10^{10} \\ 1.066 \times 10^{11} \\ 3.372 \times 10^{11} \end{array}$	$\begin{array}{c} 1.000 \times 10^5 \\ 3.162 \times 10^5 \\ 1.000 \times 10^6 \end{array}$	$\begin{array}{c} 1.334{\times}10^{4} \\ 4.218{\times}10^{4} \\ 1.334{\times}10^{5} \end{array}$	5.102 16.135 51.024	0.0510 0.1614 0.5102			

dimensionless parameters

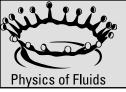
 $Ra = \frac{g\Delta\rho H^3}{\mu D}$

$$Ra_d = \frac{g\Delta\rho d^3}{\mu D}$$

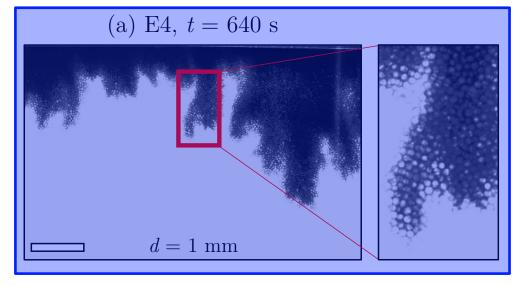
$$Pe = Ra^* Da^{1/2}$$



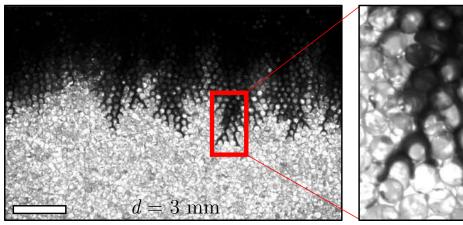
 $Re = \frac{Ra^* \, Da^{1/2}}{Sc}$



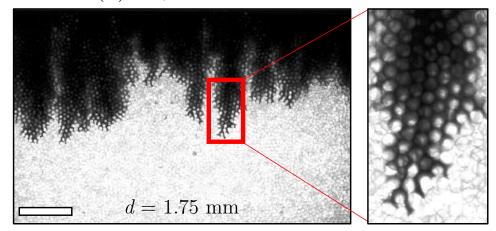
Influence of $d (\Delta \rho = 7 \text{ kg/m}^3)$



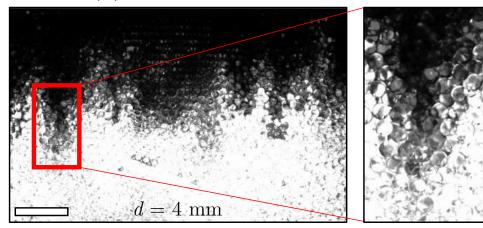
(c) E12, t = 121 s



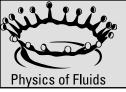
Pore-scale simulation of convective mixing in confined media *Marco De Paoli,* Physics of Fluids Group, University of Twente (b) E8, t = 206 s



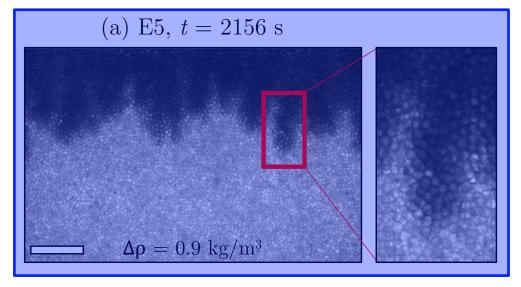
(d) E16, t = 59 s



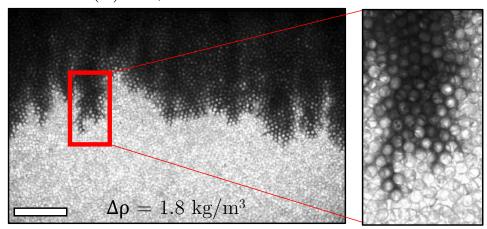




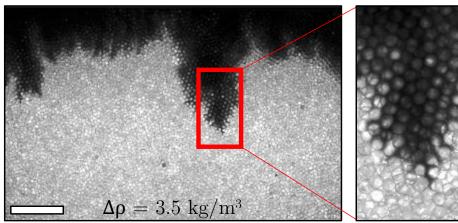
Influence of $\Delta \rho$ (*d* = 1.75 mm)



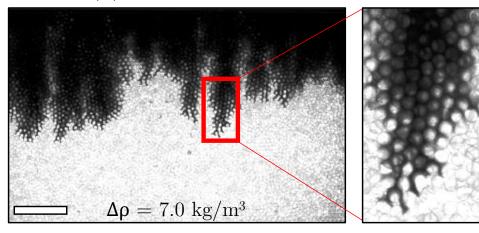
(b) E6, t = 748 s



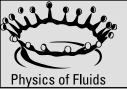
(c) E7, t = 368 s



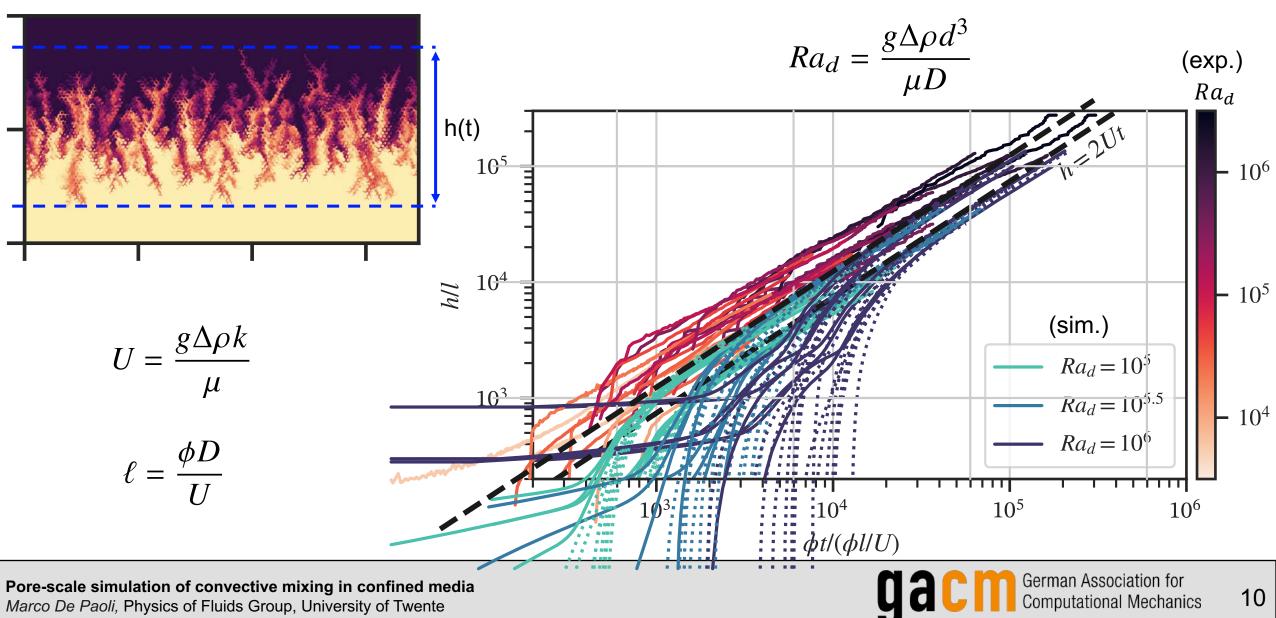
Pore-scale simulation of convective mixing in confined media *Marco De Paoli,* Physics of Fluids Group, University of Twente (d) E8, t = 206 s

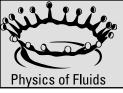




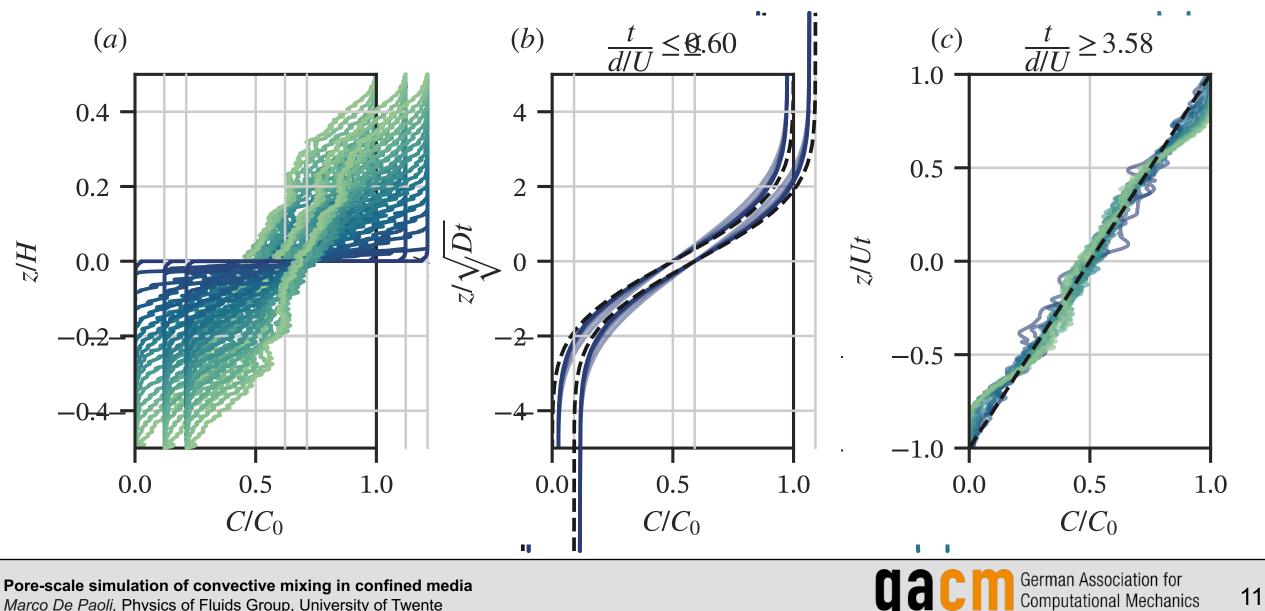


Mixing length

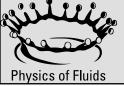


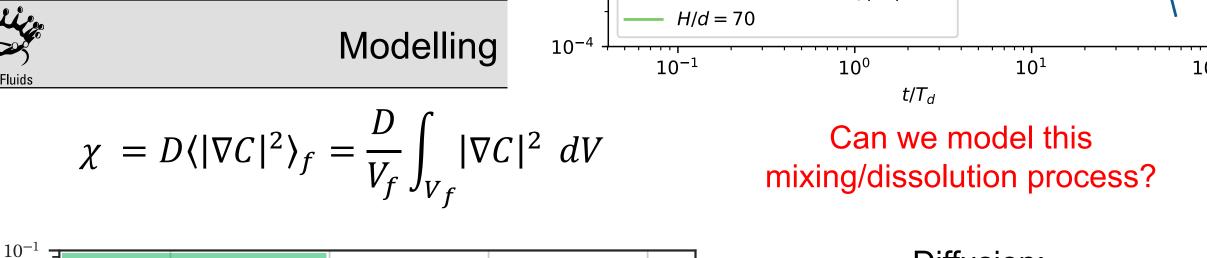


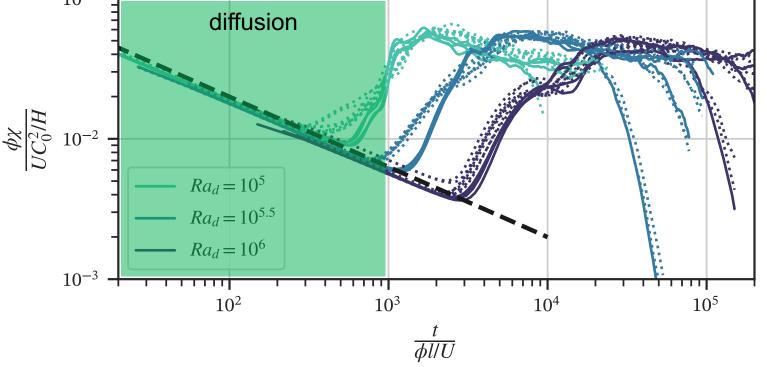
Concentration profiles



Marco De Paoli, Physics of Fluids Group, University of Twente





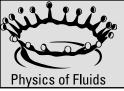


Diffusion:

$$C = C_0 + \frac{\Delta C}{2} \operatorname{erf}\left(\frac{z}{\sqrt{2\kappa t}}\right)$$
$$\partial_z C = \frac{\Delta C}{2\sqrt{\pi\kappa t}} \exp\left(-\frac{z^2}{2\kappa t}\right)$$

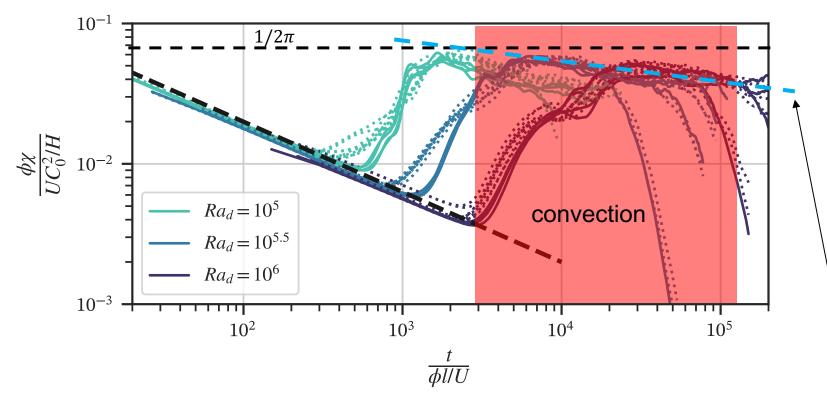
$$\chi = \kappa \langle |\nabla C|^2 \rangle = \frac{\kappa}{H} \int_{-\infty}^{\infty} |\partial_z C|^2 dz$$
$$= \sqrt{\frac{\kappa}{8\pi t}} \frac{(\Delta C)^2}{H}$$





Modelling scalar dissipation

$$\chi = D\langle |\nabla C|^2 \rangle_f = \frac{D}{V_f} \int_{V_f} |\nabla C|^2 \, dV$$



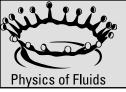
Convection $\chi = \kappa \langle |\nabla C|^2 \rangle = \kappa \frac{L_m}{H} \langle |\nabla C|^2 \rangle_{ML},$ $|\nabla C| \approx \frac{\Delta C}{2\sqrt{\pi\kappa t}}.$

 $L_m \approx 2Ut$,

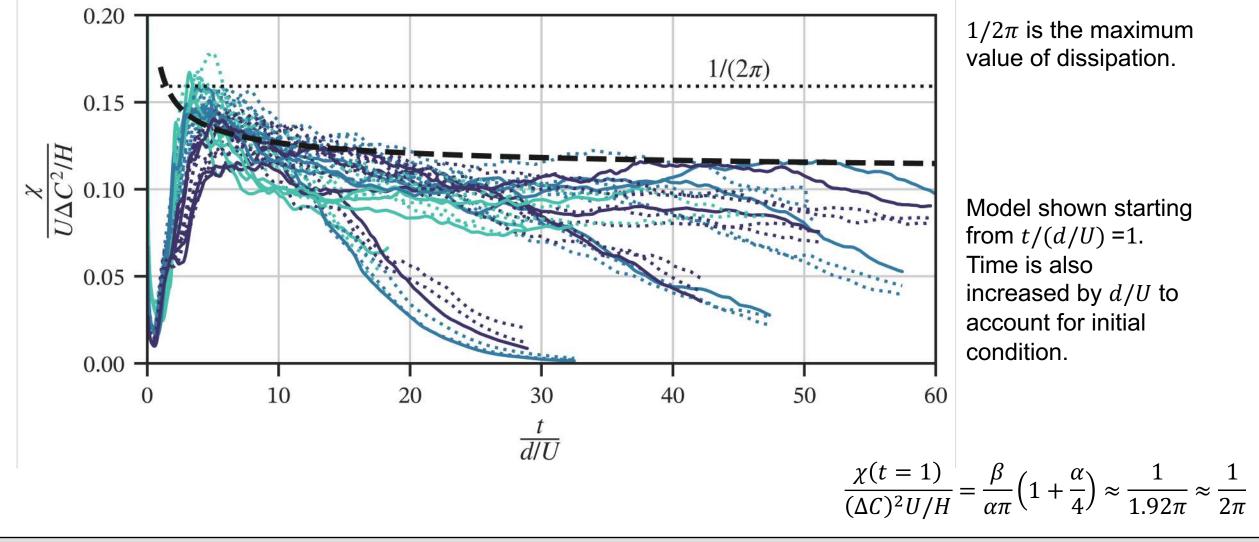
$$\chi \approx \kappa \frac{2Ut}{H} \frac{(\Delta C)^2}{4\pi\kappa t} = \frac{1}{2\pi} \frac{U_d (\Delta C)^2}{H}.$$

 $1/2\pi$ is the maximum value of dissipation. Practically, χ decreases with time

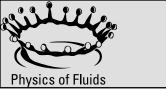


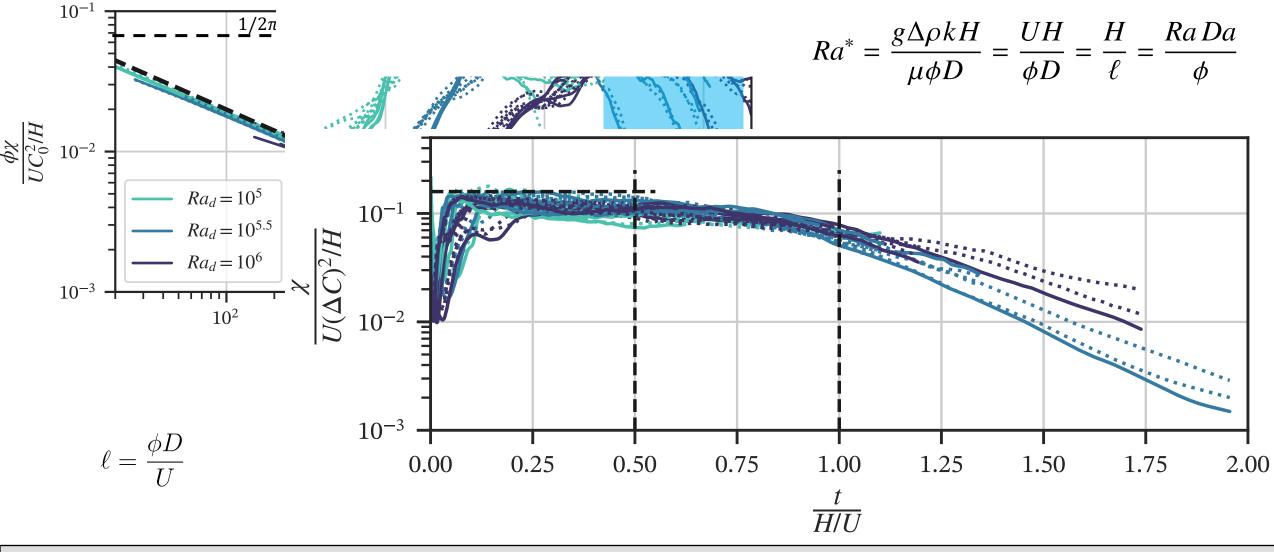


Modelling scalar dissipation

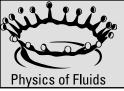








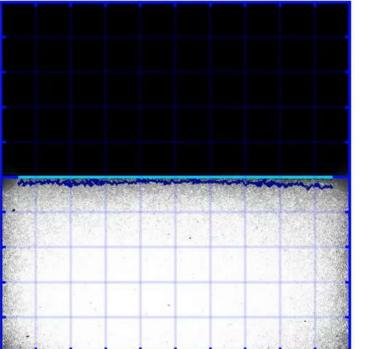




Conclusions

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- Simulations and experiments are used as complementary tools to investigate convection in porous media
- Multiple length scales are relevant at different phases of the process
- Mixing length predicted experimentally exhibits a self-similar behaviour that agrees well with theoretical prediction for convective flows in porous media
- Mixing measured numerically via mean scalar dissipation has a self-similar behaviour.
- We explain theoretically the scaling laws observed
- We plan to expand the parameters space investigated and performed simulations in three-dimensional domains



4612]

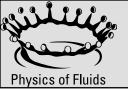
Der Wissenschaftsfonds.

Thank you for your attention! Questions?





This research was funded in part by the Austrian Science Fund (FWF) [Grant J-



High-resolution images, movies and slides are available upon request to m.depaoli@utwente.nl



17